

A sterile neutrino at MiniBooNE and IceCube

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- Gninenko's 50 MeV neutrino at LSND
- A variation of the model at MiniBooNE
- Constraints from T2K
- Implications at IceCube

Masip, Masjuan, PRD 83 (2011) 091301

Masip, Masjuan, Meloni, JHEP 01 (2013) 106

Saint Petersburg, October 2013

- Homestake, GALLEX, SAGE,... **IMB, Kamiokande, Super K, ... KEK, K2K,...**
SNO, KamLAND,... **Neutrinos have masses and mixings (!)**

$$\left\{ \begin{array}{l} \Delta m_{12}^2 \approx 7.9 \times 10^{-5} \text{ eV}^2 \\ \Delta m_{23}^2 \approx 2.5 \times 10^{-3} \text{ eV}^2 \\ \approx \Delta m_{13}^2 \end{array} \right. \quad \left\{ \begin{array}{l} \sin^2 \theta_{12} \approx 0.30 \\ \sin^2 \theta_{23} \approx 0.50 \\ \sin^2 \theta_{13} \approx 0.025 \end{array} \right.$$

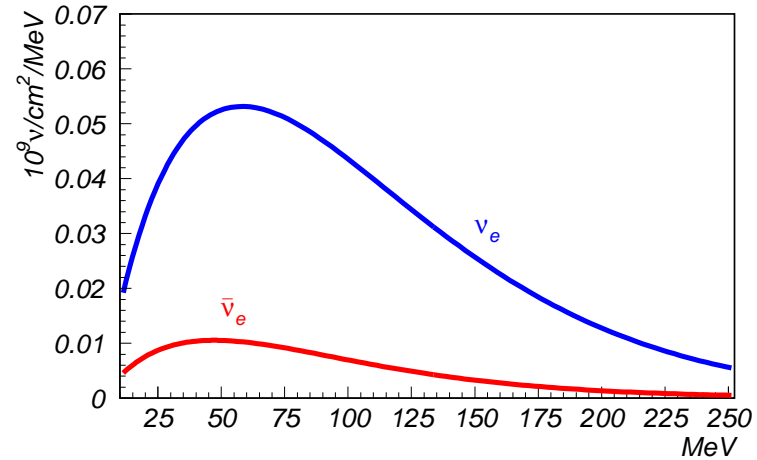
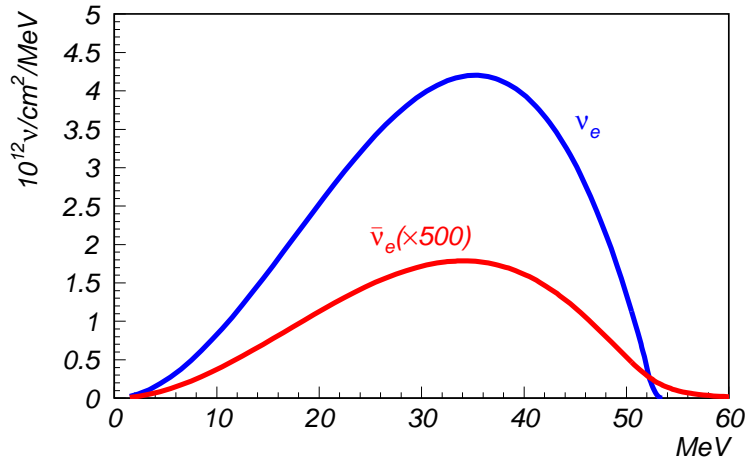
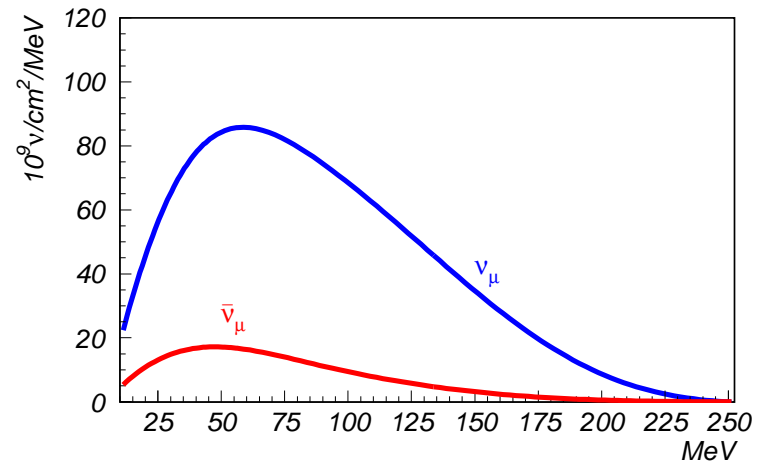
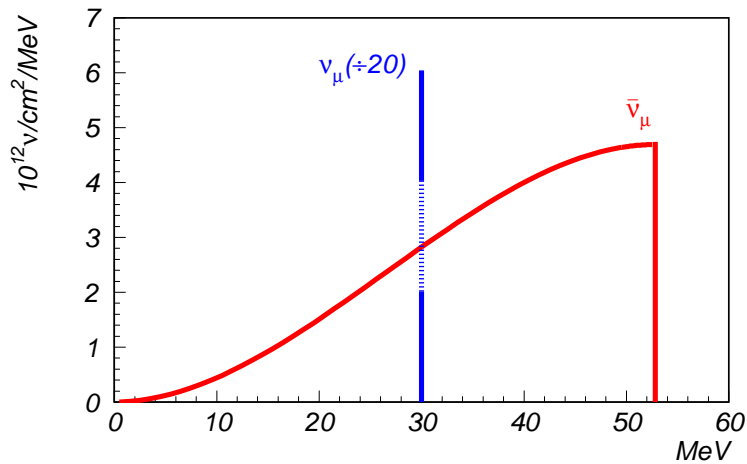
Is it $y_\nu HL\nu^c$ or $\frac{1}{\Lambda_\nu} HHLL$?

- *Persistent* anomalies in several experiments with neutrino beams from particle accelerators. **Excess of 3 events with an electron in the final state per 1000 ν_μ CC-interactions.** $\nu_\mu \rightarrow \nu_e$ oscillations inconsistent with ν -mass parameters (2 sterile neutrinos of $m \approx 1$ eV?).

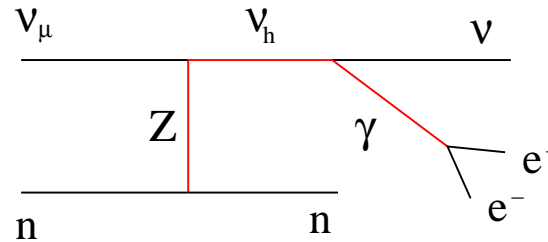
LSND, KARMEN, MiniBooNE, TRIUMF, T2K, NOMAD, IceCube

- **LSND** observed **3** electron events per **1000** $\bar{\nu}_\mu$ CC interactions. Interpreted as $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ then $\bar{\nu}_e p \rightarrow e^+ n$, with a 2.2 MeV photon from neutron capture.

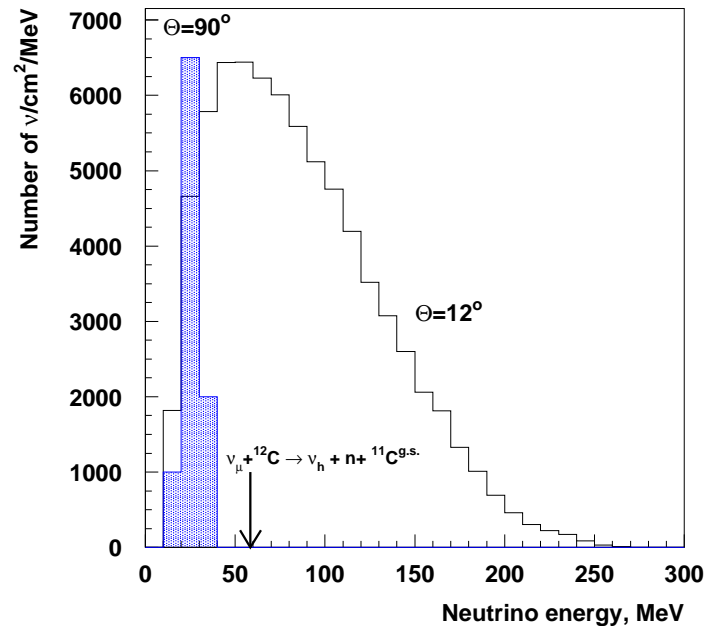
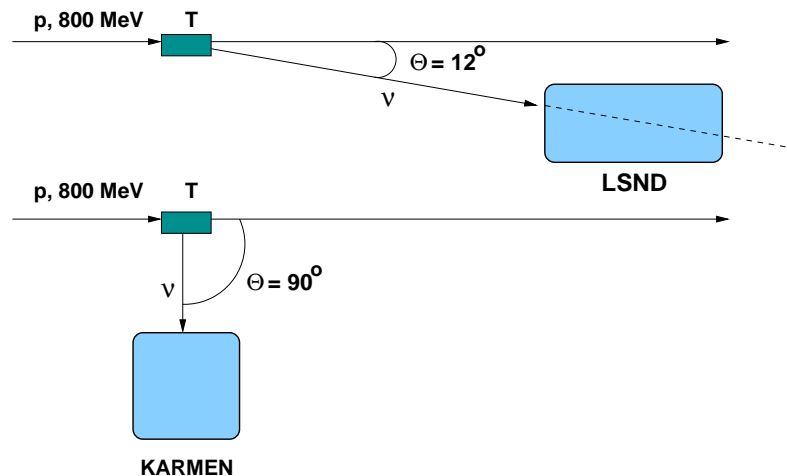
Fluxes: **DAR** (left) and **DIF** (right) $\mu^+ \rightarrow \bar{\nu}_\mu e^+ \nu_e$; $\pi^+ \rightarrow \mu^+ \nu_\mu$



Gninenko's 50 MeV neutrino hypothesis to explain LSND



- Sterile ν_h with $|U_{\mu h}|^2 \approx 10^{-3}-10^{-2}$, $\nu_h \rightarrow \nu\gamma$ with $\tau_h \lesssim 10^{-8}$ s
- KARMEN did not confirm... ν_h would be above threshold there!



ν_h would appear in up to 1% of muon and kaon decays!

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_h \rightarrow e^- \bar{\nu}_e \gamma \nu, \quad K^- \rightarrow \mu^- \bar{\nu}_h \rightarrow \mu^- \gamma \bar{\nu}$$

- Usual searches are based on decay modes with charged particles

$$\nu_h \rightarrow ee\nu, \mu e\nu, \mu\pi\nu \quad \text{not} \quad \nu_h \rightarrow \nu\gamma$$

- If ν_h is long lived ($\tau_h > 10^{-9}$ s) but light ($m_h \approx 50$ MeV), $|U_{\mu h}|^2 \approx 0.003$ does not change significantly the kinematics in μ and K decays.

- If it is short lived, muons and kaons have decay modes with photons

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu \gamma \quad 1.4 \pm 0.4\%$$

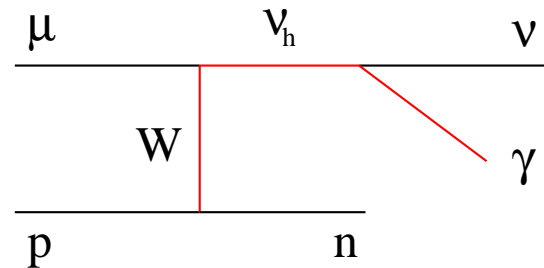
$$K^- \rightarrow \mu^- \bar{\nu}_\mu \gamma \quad 0.62 \pm 0.08\%$$

$$K^- \rightarrow \mu^- \bar{\nu}_\mu \pi^0 \quad 3.35 \pm 0.03\%$$

A recent analysis of ISTRA data seems to disfavor this possibility

- **Gninenko** proposed that ν_h can also explain the **MiniBooNE** anomaly (see below) if its lifetime is reduced to $\tau_h \leq 10^{-9}$ s (versus $\tau_h \leq 10^{-8}$ s at LSND).
- **McKeen & Pospelov** noticed that Gninenko's ν_h is inconsistent with data on muon capture plus photon at **TRIUMF**

$$R_\gamma = \frac{\Gamma_{RMC}}{\Gamma_{\text{tot}}} \Big|_{E_\gamma > 60 \text{ MeV}}$$

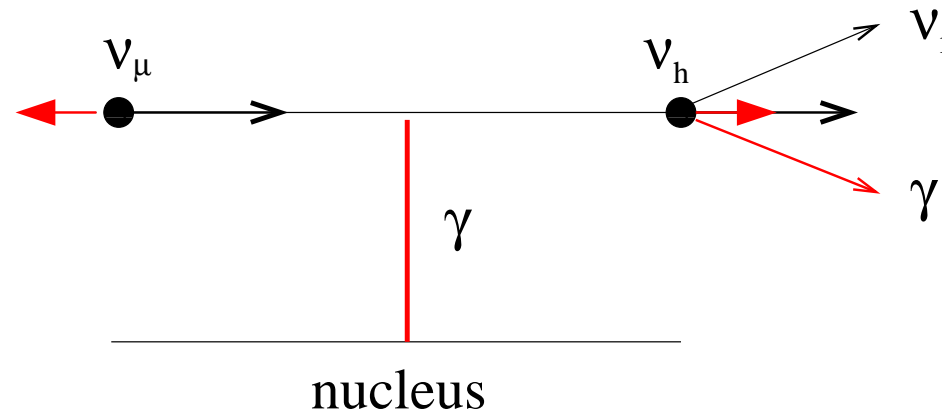


$$|U_{\mu h}|^2 \leq (2-8) \times 10^{-4} \text{ for } m_h = 40-80 \text{ MeV and } \tau_h < 10^{-9} \text{ s}$$

ν_h : factor of 3 excess **TRIUMF data: 30% excess (again a 2.5σ dev.!)**

- The cut $E_\gamma > 60$ MeV and the small size of the target volume (≈ 15 cm) make this experiment **very sensitive to the lifetime**: $\tau_h \geq 3 \times 10^{-9}$ s fits.

- Our **variation** of Gninenko's model: (i) keep a **longer lifetime**, $\tau_h \approx 5 \times 10^{-9}$ s. (ii) include ν_h production through photon exchange. (iii) ν_h a Dirac fermion.



- $\nu_h \equiv \{N_1, N_1^c\}$; N_1 mixed with ν_μ , $|U_{\mu h}|^2 \approx 0.003$; EM dipole transitions μ_{tr}^{ih} to describe ν_h production ($i = \mu$) and decay ($i = \mu, \tau, \dots$):

$$L_{eff} \subset \frac{1}{2} \mu_{tr}^{ih} \left(\bar{\nu}_h \sigma_{\mu\nu} (1 - \gamma_5) \nu_i + \bar{\nu}_i \sigma_{\mu\nu} (1 + \gamma_5) \nu_h \right) \partial^\mu A^\nu$$

$$\tau_h = 5 \times 10^{-9} \text{ s implies } \sqrt{\sum_i (\mu_{tr}^{ih})^2} = 7 \times 10^{-6} \text{ GeV}^{-1} = 2 \times 10^{-8} \mu_B$$

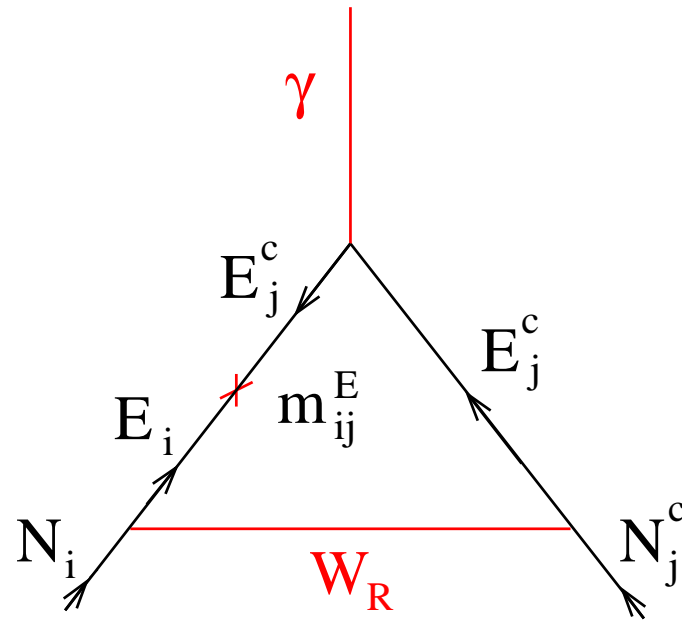
$$\text{MiniBooNE will require } \mu_{tr}^{\mu h} = 2 \times 10^{-9} \mu_B$$

$$SU(2)_R \times U(1)_{B-L} \rightarrow U(1)_Y : \chi = \begin{pmatrix} \langle \chi^0 \rangle \\ \chi^- \end{pmatrix} \quad \chi^c = \begin{pmatrix} \chi^{c+} \\ \langle \chi^{c0} \rangle \end{pmatrix}$$

$$Y = T_R^3 + B-L$$

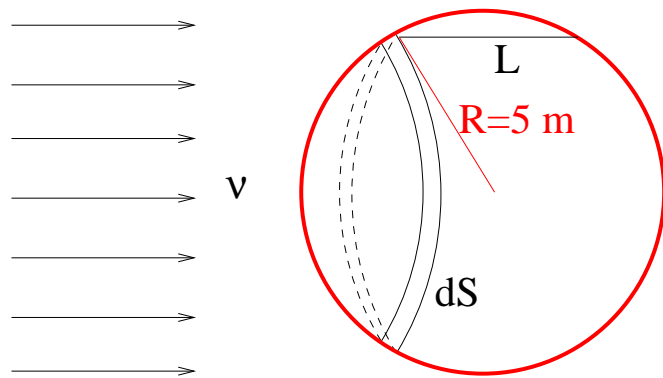
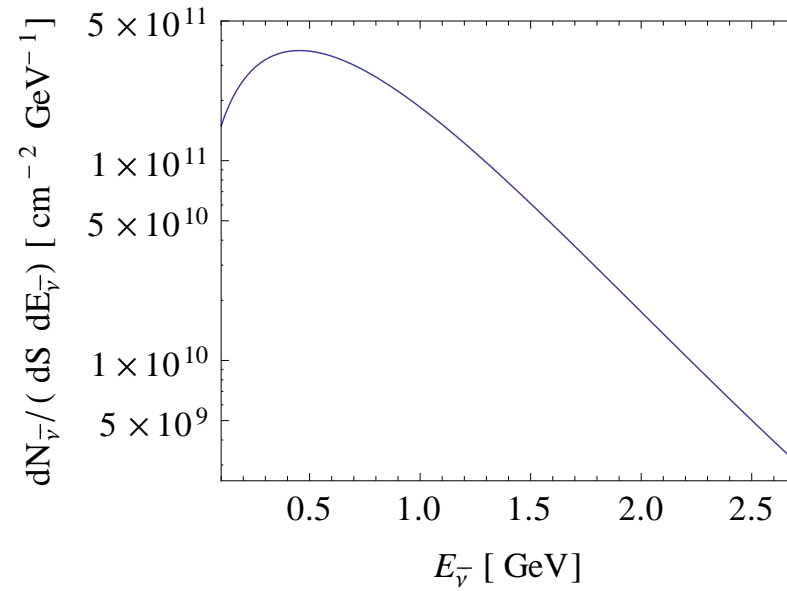
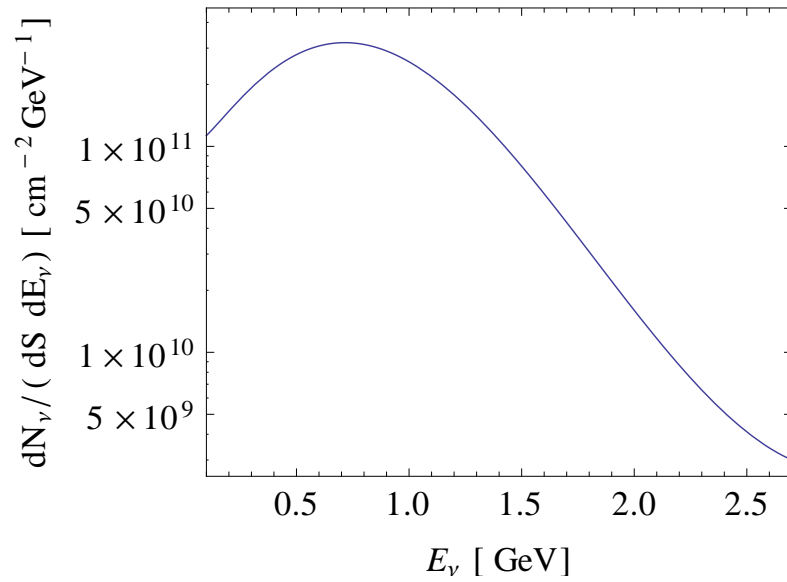
$$Q = T_L^3 + Y$$

$$V_\mu' = c_\alpha V_\mu + s_\alpha N_1$$



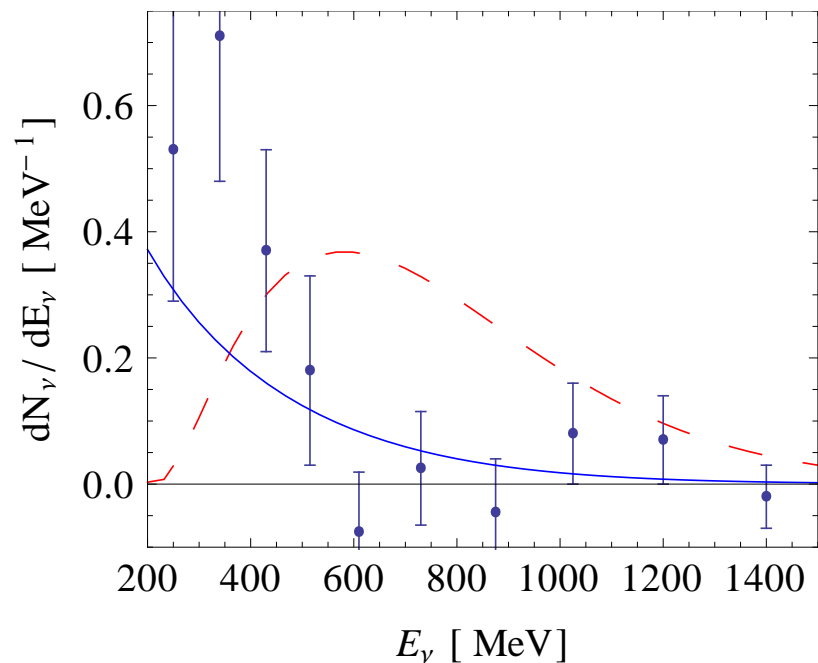
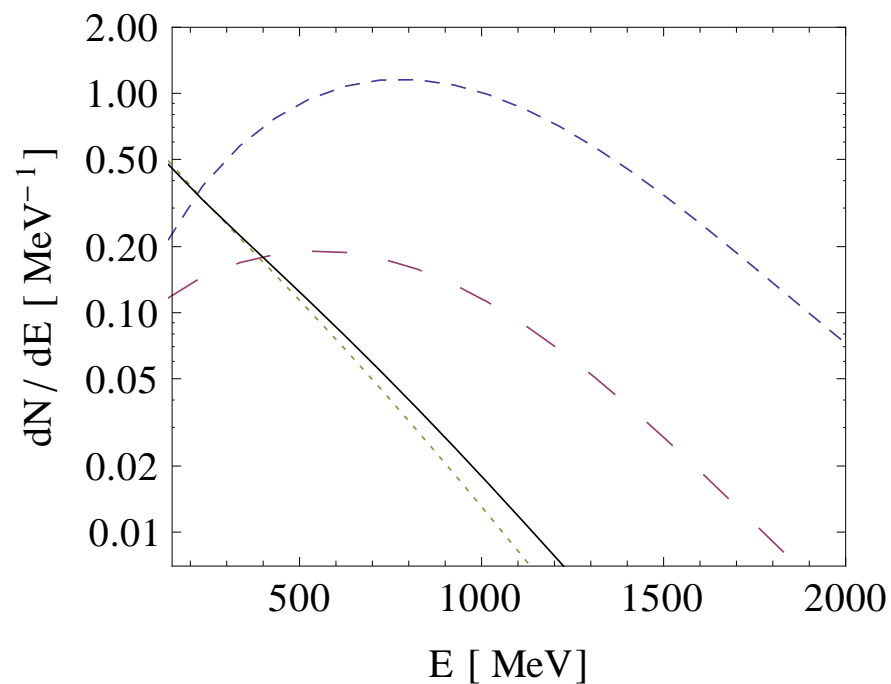
$$L_1 = \begin{pmatrix} N_1 \\ E_1 \end{pmatrix} \quad L_1^c = \begin{pmatrix} E_1^c \\ N_1^c \end{pmatrix}$$

$$W \subset m_h L_1 L_1^c + \frac{1}{\Lambda_1} (L_1^c \chi^c) (L_1 \chi)$$

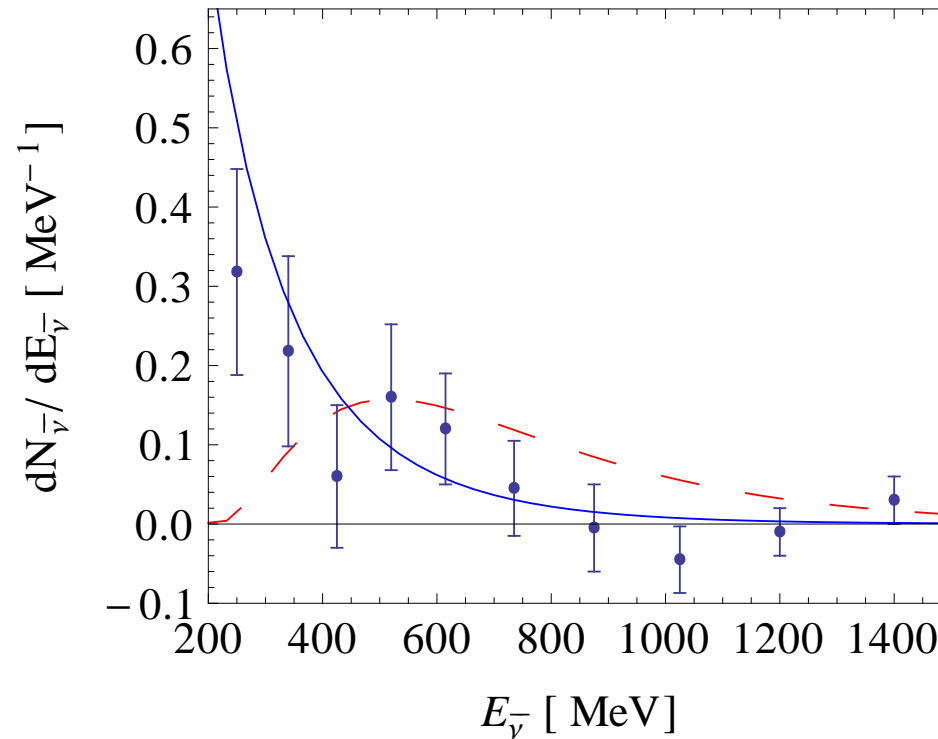


$$p_i = \frac{\sigma \rho L}{m_{CH_2}}$$

$$p_{i+d} = \frac{\sigma \rho L}{m_{CH_2}} \left(1 - \frac{\lambda_d}{L} \left(1 - e^{-\frac{L}{\lambda_d}} \right) \right)$$



Left: Energy distribution of ν_h produced in the detector (dashes), of ν_h decaying inside the detector (long dashes), of photons from $\nu_h \rightarrow \nu_i \gamma$ (dots), and of ν_h events reconstructed as CC interactions (solid). **Right:** Energy distribution of ν_h events reconstructed as CC interactions (solid), of events from neutrino oscillations for $\sin^2(2\theta) = 0.004$ and $\Delta m^2 = 1 \text{ eV}^2$ (long dashes), and excess at MiniBooNE in the neutrino mode (5.58×10^{20} POT)



Antineutrino mode (11.27×10^{20} POT)

- The decay length ($\lambda_{dec} > R$) and the helicity (+ for ν_h , - for $\bar{\nu}_h$) imply that the MiniBooNE excess concentrates at low energies, just as it is observed.

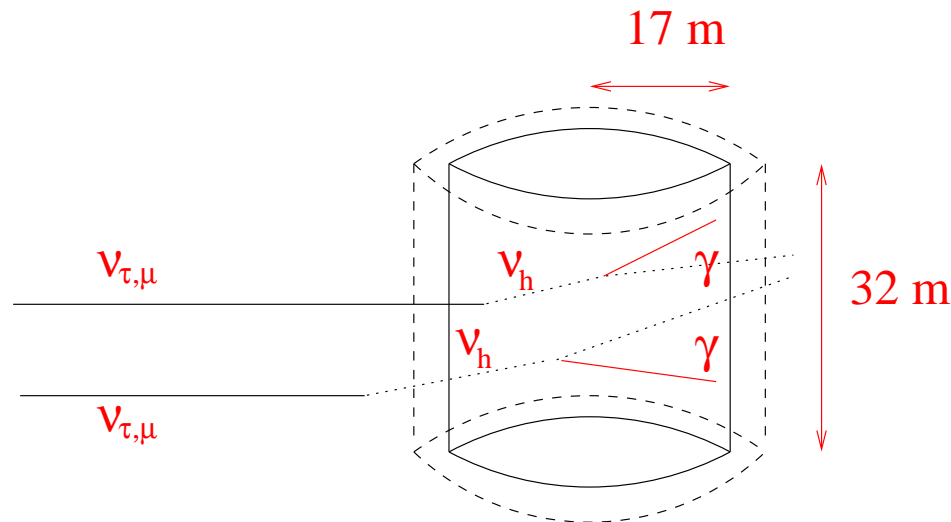
If $\text{BR}(\nu_h \rightarrow \nu_\mu \gamma) \approx 1\%$, is $\text{BR}(\nu_h \rightarrow \nu_\tau \gamma) \approx 99\%$?

At T2K we expect some ν_e ,

$$P_{\nu_\mu \rightarrow \nu_e} \approx \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \frac{1.27 \Delta m_{23}^2 (\text{eV}^2) L (\text{km})}{E_\nu (\text{GeV})}$$

but most neutrinos are ν_τ ,

$$P_{\nu_\mu \rightarrow \nu_\tau} \approx \sin^2 2\theta_{23} \sin^2 \frac{1.27 \Delta m_{23}^2 (\text{eV}^2) L (\text{km})}{E_\nu (\text{GeV})}$$

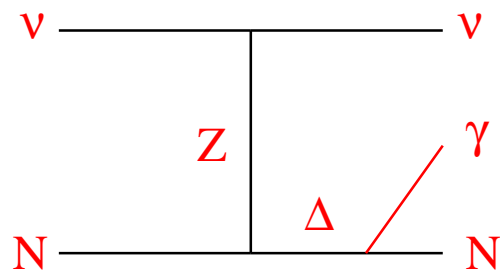


- $\nu_\mu \rightarrow \nu_e$ oscillations with $\sin^2 2\theta_{13} = 0.1$: **6 events**
- $\nu_\mu Z \rightarrow \nu_h Z$: **1.1 events** (75% from ν_h produced outside).
- ν_h must decay 99% of the times into another sterile neutrino!

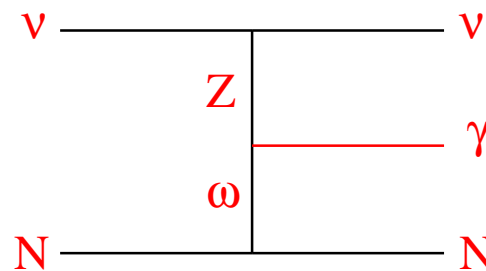
$$\mu_{\text{tr}}^{\tau h} < \mu_{\text{tr}}^{\mu h} \quad \text{and} \quad \text{BR}(\nu_h \rightarrow \nu_{h'} \gamma) = 0.99$$

- Initial events seem to be distributed near the point of entrance into the detector. $\nu_\tau Z \rightarrow \nu_h Z$ events could explain that: When ν_h is produced outside the detector $\lambda_{\text{dec}} \approx d$
- The tracking system in the near detector (ND280) can distinguish electrons from photons: we expect **3 ν_h events per 1000 ν_μ CC interactions**

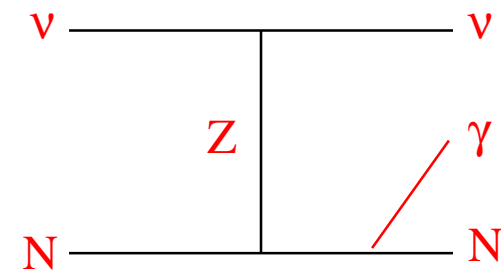
- **MicroBooNE** will investigate whether the low-energy excess at MiniBooNE is caused by **electron** or by **photon** events,



Delta resonance



pseudo Chern-Simons



Compton

- There are observables that may distinguish this photon background from the $\nu_h \rightarrow \gamma \nu_{h'}$ hypothesis:

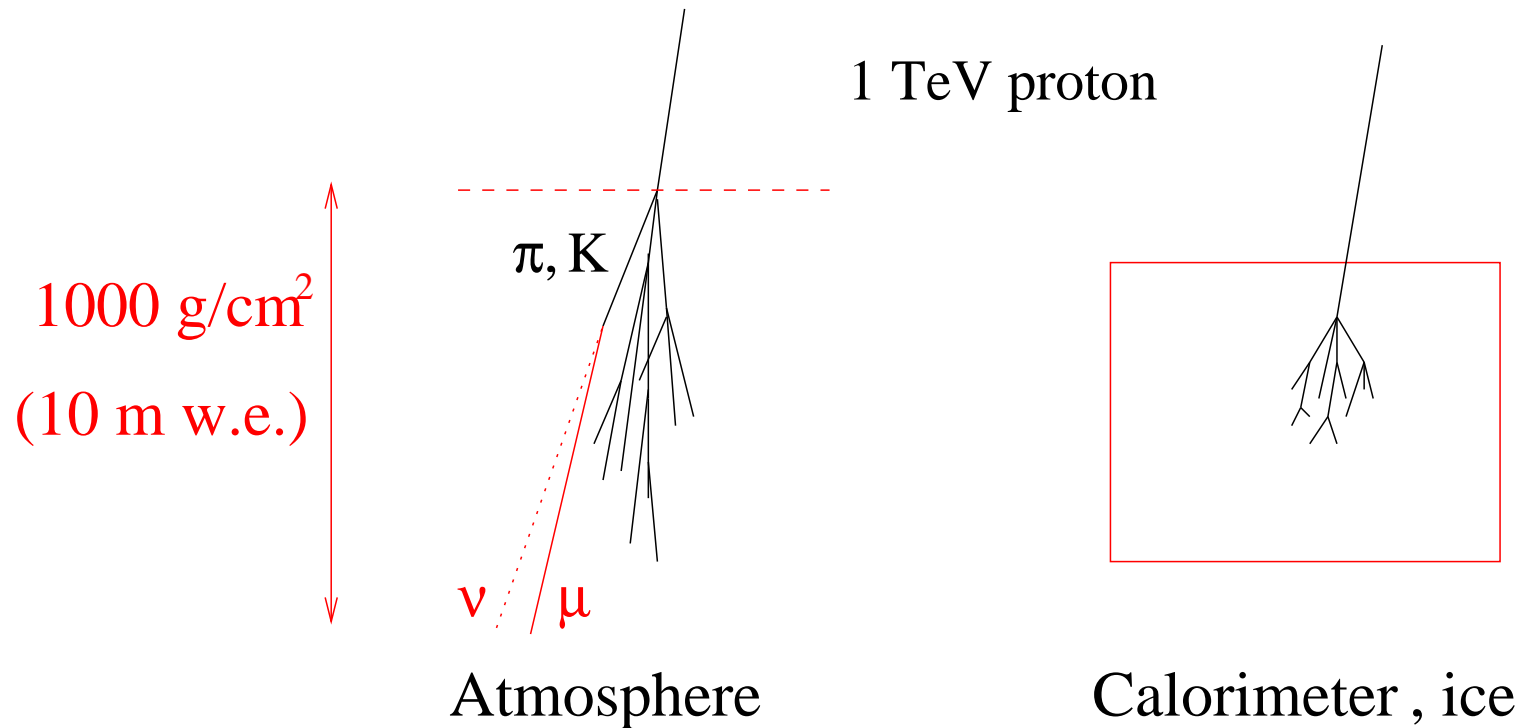
The event distribution inside the detector is flat for the background events,

but $\propto (1 - e^{-z/\lambda_d}) \approx \frac{z}{\lambda_d}$ for heavy neutrino events.

LSND: anomaly at $L \approx 30$ m for $E \approx 40$ MeV

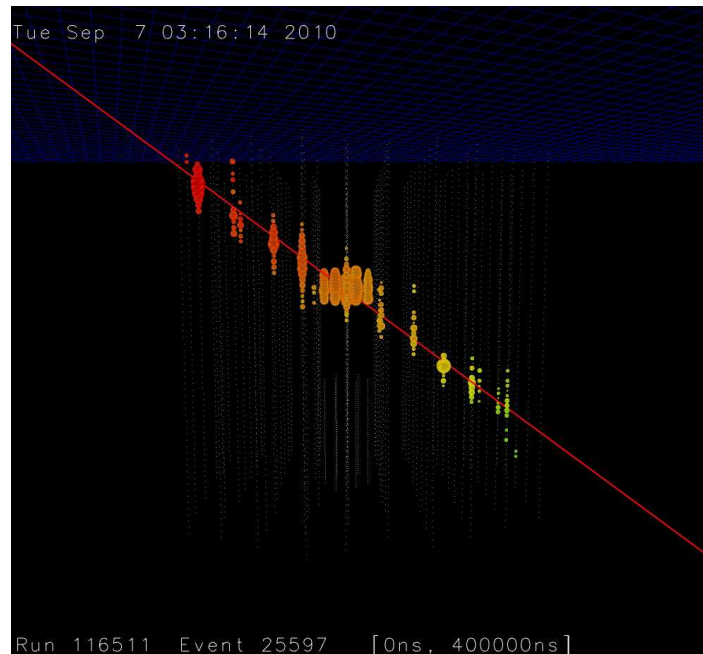
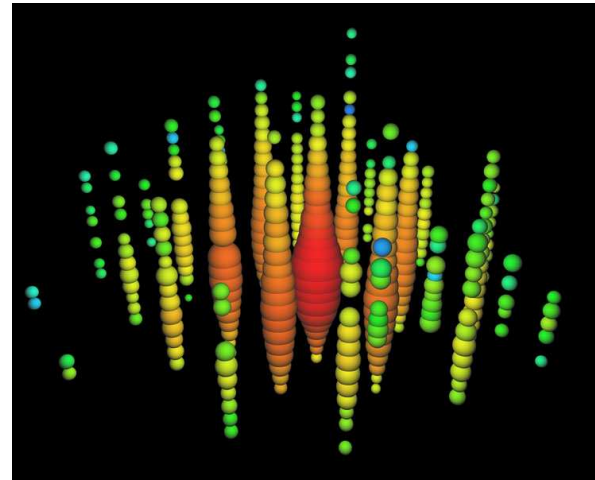
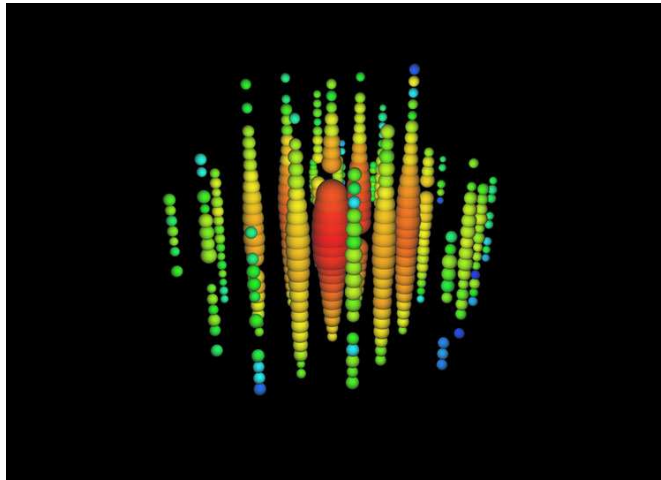
MiniBooNE: anomaly at $L \approx 500$ m for $E \approx 300$ MeV

ν telescopes: anomaly at $L \approx 10\text{--}100$ km for $E \approx 1$ TeV ??



- Atmospheric muons and neutrinos can be seen at IceCube

- Bert, Ernie & Muon



- Analytical meson and lepton fluxes [*Z–moment* method, Gaisser, Lipari]

Set of coupled differential equations that describe the evolution with the atmospheric depth t (in g/cm^2) of the fluxes of *parent* hadrons (ϕ_H with $H = p, n, \pi^\pm, K^\pm, K_L$) and of any particles that may result from their decay or their collision with an air nucleus: $[\phi_H(E, \theta, t)]$

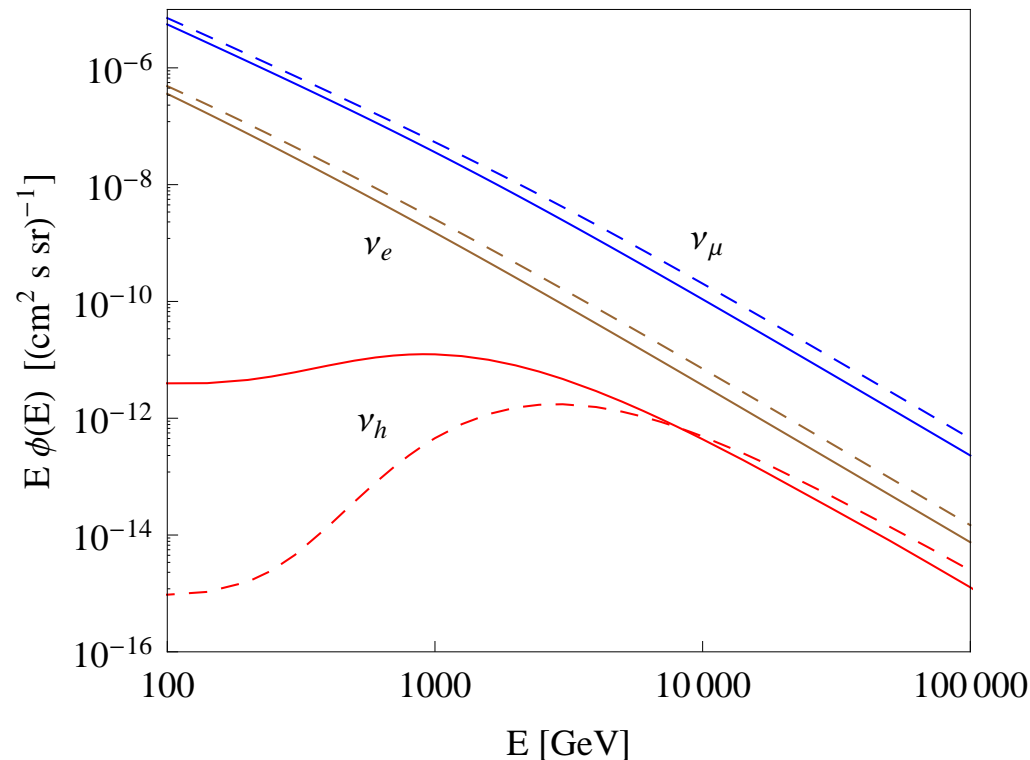
$$\frac{\partial \phi_H}{\partial t} = -\frac{\phi_H}{\lambda_{\text{dec}}^H} - \frac{\phi_H}{\lambda_{\text{int}}^H} + \sum_{H'} S_{H'H}$$

Source: $S_{H'H} = \frac{\phi_{H'}}{\lambda_{\text{int}}^{H'}} Z_{H'H}$ Z–factors: $Z_{H'H} = \int_0^1 dx x^{\alpha-1} F_{H'H}$

$F_{H'H}(x)$: distribution of the fraction of energy taken by H after a H' –air collision ($x = E_H/E_{H'}$). Primary all nucleon flux: $\Phi_N \propto E^{-\alpha}$

$$B(K^+ \rightarrow \mu^+ \nu_h) \approx B(K^+ \rightarrow \mu^+ \nu) \times |U_{\mu h}|^2 \bar{\rho}_h \quad \bar{\rho}_h \approx 1 + \frac{m_h^2}{m_\mu^2}$$

- Neutrino fluxes ($\nu_i + \bar{\nu}_i$) at sea level for $\theta = 0$ (solid) and $\theta = 60^\circ$ (dashes)



Z-moment method

$$m_h = 60 \text{ MeV}$$

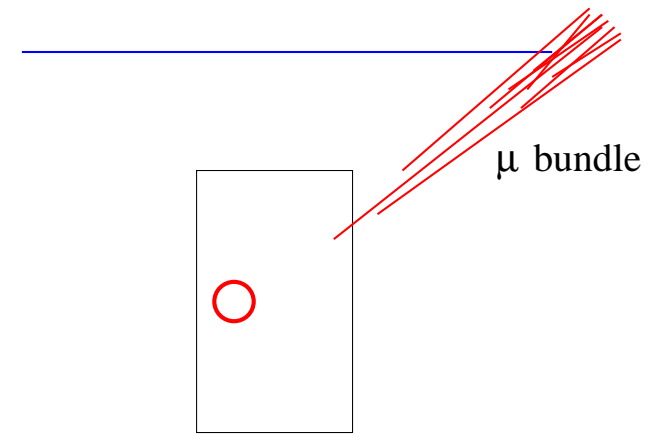
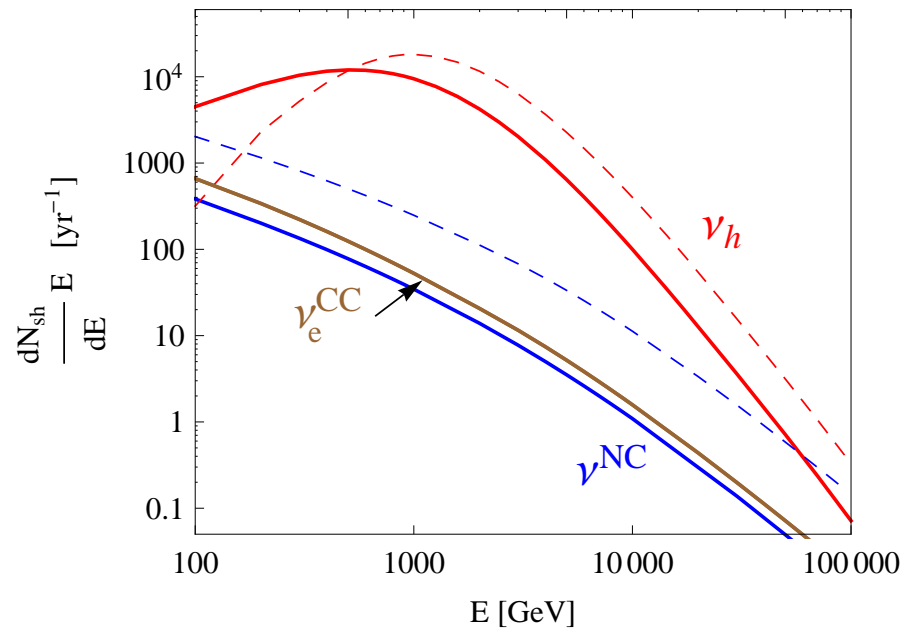
$$|U_{\mu h}|^2 = 0.005$$

$$\tau_h = 10^{-9} \text{ s}$$

$$\lambda_{dec} = 5 \text{ km at } E = 1 \text{ TeV}$$

[PRD83(2011)091301]

- Contained events at ANTARES and the DeepCore in IceCube. In dashes the energy distribution of the parent neutrino.

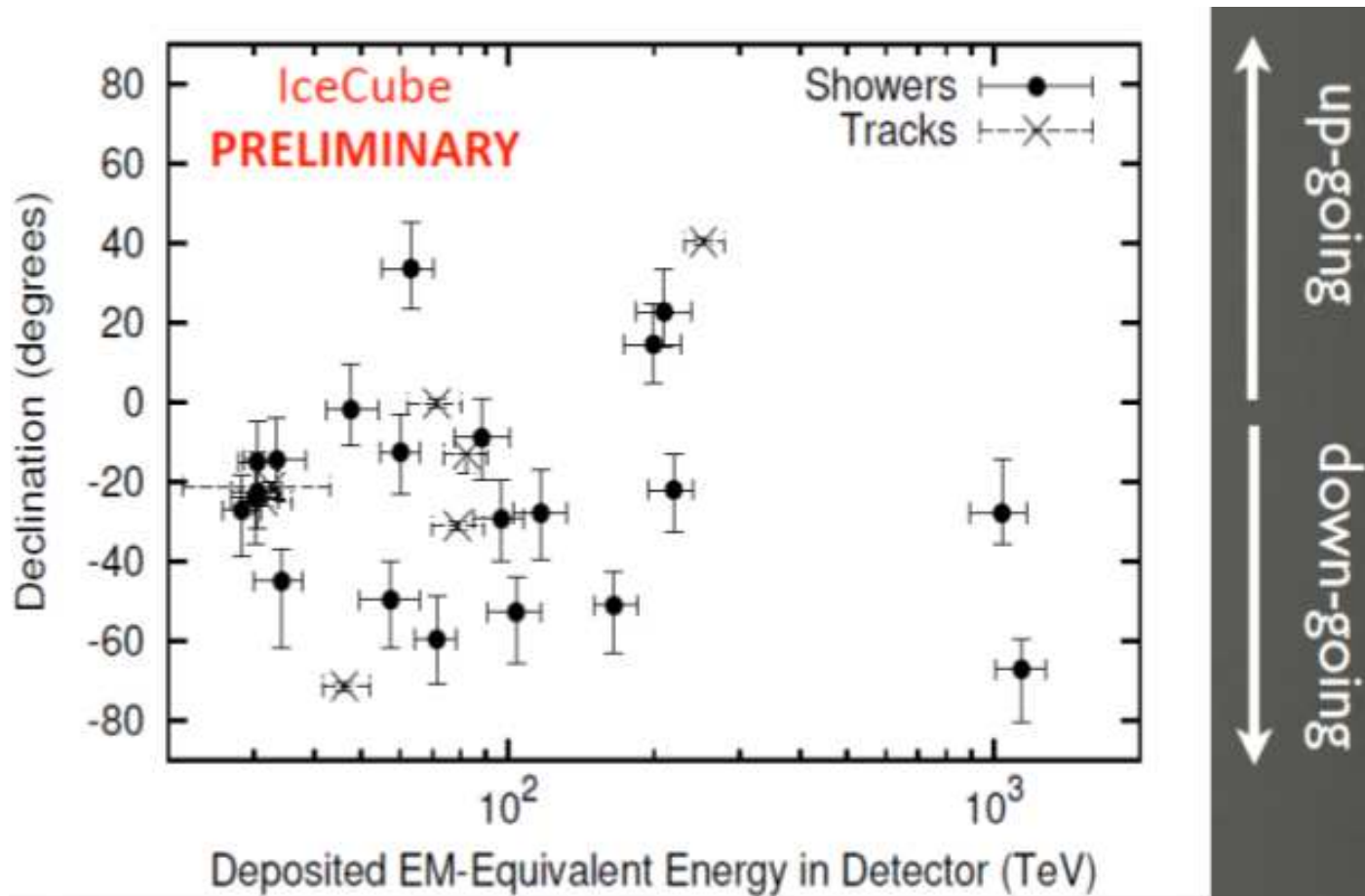


- 14000 $\nu_h \rightarrow \gamma\nu$ events of energy above 500 GeV per year, versus 220 standard events ($\nu_e N \rightarrow eX$ and $\nu_{\mu,e} N \rightarrow \nu_{\mu,e} X$)

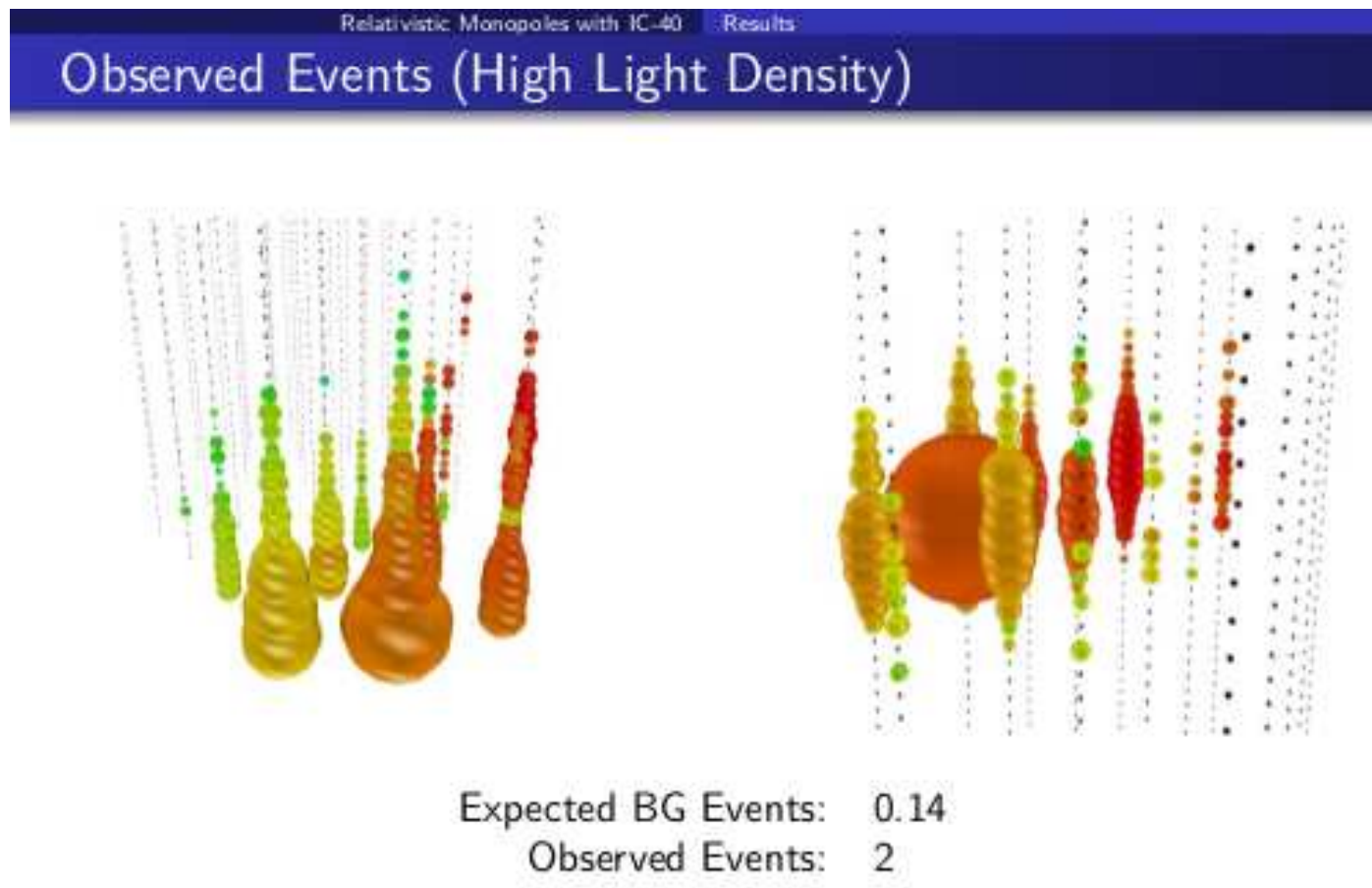
- At energies below 100 GeV ν_h does not reach the telescope, above 1000 TeV its decay length becomes too large and the signal vanishes.

- Neutrinos ν_h produced in the atmosphere and decaying inside IceCube would produce an **excess of contained events** (similar to ν_e CC interactions or inelastic NC collisions) at energies 1–1000 TeV
- This excess would only appear in **downgoing or near-horizontal events** (no ν_h upgoing events)
- Most of these events (specially the ones from small zenith angles) would be contaminated with muons. **Excess of muons plus contained cascade**

recent data from IceCube at IPA 2013 ...



... from an IceCube presentation at *Exotics with neutrino telescopes 2013* ...



SUMMARY

- Neutrino physics has progressed a great deal during the past 20 years, but (i) basic questions are still unanswered and (ii) some *persistent* anomalies should be clarified (MicroBooNE next year?)
- A 50 MeV neutrino ν_h mixed with the muon flavor ($|U_{\mu h}|^2 \approx 0.003$) with a lifetime $c\tau = 1.5$ m, produced $\nu_\mu Z \rightarrow \nu_h Z$ and decaying $\nu_h \rightarrow \nu_{h'} \gamma$ through **electromagnetic dipole transitions** could explain LSND, KARMEN, TRIUMF, MiniBooNE.
- An excess of **contained events** at IceCube could be correlated with the LSND and MiniBooNE anomalies. These events would only be **downgoing and quasi-horizontal**, possibly *contaminated* by muons from the parent air shower.