

# ***Factorization of R-matrix, Baxter Q-operators and SOV***

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## motivation: Separation of Variables

Let  $\Psi(\mathbf{z}) = \Psi(z_1, \dots, z_N)$  is eigenfunction of a quantum integrable system

Expand  $\Psi(z_1, \dots, z_N)$  into a basis of eigenfunctions of another(auxiliary) integrable system

$$\Psi(\mathbf{z}) = \int U(\mathbf{z}, \mathbf{x}) Q(\mathbf{x}) \mu(\mathbf{x}) d^N \mathbf{x}$$

Separation of variables  $\leftrightarrow$  factorization of  $Q(\mathbf{x})$ :  $Q(\mathbf{x}) = Q(x_1)Q(x_2) \cdots Q(x_N)$

✓  $U(\mathbf{z}, \mathbf{x})$  can be constructed iteratively and  $Q(x)$  is solution of Baxter equation

[Sklyanin]

main example

✓ integrable system: XXX spin chain  $S_k = z_k \partial_k + \ell$  ;  $S_k^+ = z_k^2 \partial_k + 2\ell z_k$  ;  $S_k^- = -\partial_k$

$$L_k(u) = \begin{pmatrix} u + S_k & S_k^- \\ S_k^+ & u - S_k \end{pmatrix} ; T(u) = L_1(u) \cdots L_N(u) = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix}$$

$$t(u) = \text{tr } T(u) = A(u) + D(u) = 2 \cdot u^N + u^{N-2} \cdot \mathbf{q}_2 + u^{N-3} \cdot \mathbf{q}_3 + \cdots + \mathbf{q}_N ; [t(u), t(v)] = 0$$

$$\mathbf{q}_k \Psi(z_1, \dots, z_N) = q_k \Psi(z_1, \dots, z_N)$$

✓ auxiliary integrable system  $[B(u), B(v)] = 0$  ;  $\mathbf{x} = x_1, \dots, x_{N-1}$

$$B(u)U(\mathbf{z}, \mathbf{x}, p) = p(u - x_1) \cdots (u - x_{N-1}) \cdot U(\mathbf{z}, \mathbf{x}, p) ; S^- U(\mathbf{z}, \mathbf{x}, p) = p \cdot U(\mathbf{z}, \mathbf{x}, p)$$

✓ Baxter equation

$$t(x)Q(x) = \Delta_+(x) \cdot Q(x+1) + \Delta_-(x) \cdot Q(x-1)$$

## examples and plan

$Q$ -operators  $\leftrightarrow$  R.Baxter

Baxter, Bazhanov, Stroganov, Lukyanov, Zamolodchikov, Volkov, Faddeev, Kashaev, Pasquier, Gaudin, Pronko, Smirnov, Korff, Sklyanin, Kuznetsov, Salerno, SD, Korchemsky, Manashov, Kulish, Kharchev, Lebedev, Gerasimov, Oblezin, Bytsko, Teschner ...

SOV:

- ✓ quantum Toda chain: Gutzwiller, Sklyanin, Pasquier-Gaudin, Kharchev, Lebedev, Gerasimov, Oblezin, Silantyev
- ✓ XXX spin chain: Faddeev, Korchemsky, SD, Manashov, Lipatov
- ✓ q-deformed XXX spin chain: Manashov, Kirch, Bytsko, Teschner
- ✓  $SL(n)$ -invariant spin chain: Sklyanin, Smirnov

Plan: algebraic part – demonstration of the connection:  $R$  - matrix  $\leftrightarrow U(\mathbf{z}, \mathbf{x}, p) \leftrightarrow Q$ -operators

- ✓ principal series representation of  $SL(2, C)$
- ✓ solution of general Yang-Baxter equation and its factorization
- ✓ factorization of general transfer-matrix into  $Q$ -operators
- ✓ iterative diagonalization of  $B(u)$
- ✓ repetition of all steps except last ones for  $SL(n, C)$

[SD,Korchemsky,Manashov]

[SD,Manashov]

## representations of $GL(2, C)$

$$z = \begin{pmatrix} 1 & 0 \\ z & 1 \end{pmatrix} ; \quad h = \begin{pmatrix} h_{11} & h_{12} \\ 0 & h_{22} \end{pmatrix} ; \quad g = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

**main construction:** representation with labels  $(\sigma_1, \sigma_2)$  and  $(\bar{\sigma}_1, \bar{\sigma}_2)$ ,  $\sigma_k - \bar{\sigma}_k \in \mathbb{Z}$

$$g^{-1} \cdot z = z' \cdot h \quad ; \quad T(g) \Phi(z) = [h_{11}]^{\sigma_1+1} \cdot [h_{22}]^{\sigma_2+2} \cdot \Phi(z')$$

$[z]^a \equiv z^a \cdot \bar{z}^{\bar{a}} = |z|^{2a} \cdot \bar{z}^{\bar{a}-a}$ , where  $z$  and  $\bar{z}$  are complex conjugate, but  $a - \bar{a} \in \mathbb{Z}$

**explicit formulae:**

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix}^{-1} \begin{pmatrix} 1 & 0 \\ z & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ \frac{-c+az}{d-bz} & 1 \end{pmatrix} \begin{pmatrix} \frac{d-bz}{ad-bc} & -\frac{b}{ad-bc} \\ 0 & \frac{1}{d-bz} \end{pmatrix}$$

$$T(g) \Phi(z, \bar{z}) = [ad - bc]^{1-\sigma_1} \cdot [d - bz]^{\sigma_1-\sigma_2-1} \cdot \Phi \left( \frac{-c + az}{d - bz}, \frac{-\bar{c} + \bar{a}\bar{z}}{\bar{d} - \bar{b}\bar{z}} \right)$$

**intertwining operator:** representations  $(\sigma_1, \sigma_2)$  and  $(\sigma_2, \sigma_1)$  are equivalent

$$S \cdot T^{(\sigma_1, \sigma_2)}(g) = T^{(\sigma_2, \sigma_1)}(g) \cdot S \quad ; \quad S = (i\partial_z)^{\sigma_1-\sigma_2} (i\partial_{\bar{z}})^{\bar{\sigma}_1-\bar{\sigma}_2} = [i\partial_z]^{\sigma_1-\sigma_2}$$

## representations of $GL(2, C)$

generators:  $g = 1 + \epsilon \cdot \mathbf{e}_{ik}$

$$T(g) \Phi(z, \bar{z}) = \Phi(z, \bar{z}) + (\epsilon \cdot E_{ik} + \bar{\epsilon} \cdot \bar{E}_{ik}) \Phi(z, \bar{z}) + \dots$$

$$E_{11} = z\partial_z + 1 - \sigma_1 ; E_{12} = z^2\partial_z + (1 - \sigma_1 + \sigma_2)z ; E_{21} = -\partial_z ; E_{22} = -z\partial_z - \sigma_2$$

commutation relations

$$[E_{ik}, E_{nm}] = \delta_{kn} E_{im} - \delta_{im} E_{nk}$$

L-operator

$$L(u) = \begin{pmatrix} u + E_{11} & E_{21} \\ E_{12} & u + E_{22} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ z & 1 \end{pmatrix} \begin{pmatrix} u - \sigma_1 & -\partial_z \\ 0 & u - \sigma_2 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -z & 1 \end{pmatrix}$$

natural combinations  $u_1 = u - \sigma_1 ; u_2 = u - \sigma_2$

intertwining operator exchanges parameters in L-operator

$$S = [i\partial_z]^{\sigma_{12}} = [i\partial_z]^{u_2 - u_1} ; [i\partial_z]^{u_2 - u_1} \cdot L(u_1, u_2) = L(u_2, u_1) \cdot [i\partial_z]^{u_2 - u_1}$$

## principal series representations of $SL(2, C)$

$ad - bc = 1 \rightarrow$  only differences  $\sigma_{12} = \sigma_1 - \sigma_2$  and  $\bar{\sigma}_{12} = \bar{\sigma}_1 - \bar{\sigma}_2$  enter

$$T(g) \Phi(z, \bar{z}) = [d - bz]^{\sigma_{12}-1} \cdot \Phi\left(\frac{-c + az}{d - bz}, \frac{-\bar{c} + \bar{a}\bar{z}}{\bar{d} - \bar{b}\bar{z}}\right)$$

**unitary representations:** representation with labels  $\sigma_{12} = -\frac{n}{2} + i\lambda$  ;  $\bar{\sigma}_{12} = \frac{n}{2} + i\lambda$  ;  $n \in \mathbb{Z}$  ,  $\lambda \in \mathbb{R}$

$$\langle \Phi_1 | \Phi_2 \rangle = \int d^2 z \overline{\Phi_1(z, \bar{z})} \Phi_2(z, \bar{z}) \ ; \ \langle T(g)\Phi_1 | T(g)\Phi_2 \rangle = \langle \Phi_1 | \Phi_2 \rangle$$

**intertwining operator is unitary**

$$S = (i\partial_z)^{-\frac{n}{2}+i\lambda} (i\partial_{\bar{z}})^{\frac{n}{2}+i\lambda} \rightarrow S^\dagger = (i\partial_{\bar{z}})^{-\frac{n}{2}-i\lambda} (i\partial_z)^{\frac{n}{2}-i\lambda} = S^{-1}$$

**symmetric labels:** labels  $(\sigma_1, \sigma_2)$  and  $(\sigma_1 + \sigma, \sigma_2 + \sigma)$  correspond the same representation of  $SL(2, C)$ . We fix the freedom by condition  $\sigma_1 + \sigma_2 = 1$ .

# Yang-Baxter equation

general Yang-Baxter equation

$$R_{12}(u-v)R_{13}(u)R_{23}(v) = R_{23}(v)R_{13}(u)R_{12}(u-v)$$

defining relation

$$R_{12}(u-v)L_1(u)L_2(v) = L_2(v)L_1(u)R_{12}(u-v)$$

it is useful to extract permutation:  $R_{12} = P_{12} \check{R}_{12}$  ,  $P_{12} \Phi(z_1, z_2) = \Phi(z_2, z_1)$

$$\check{R}(u-v) L_1(u_1, u_2) L_2(v_1, v_2) = L_1(v_1, v_2) L_2(u_1, u_2) \check{R}(u-v)$$

permutation group  $\mathfrak{S}_4$  of four parameters:  $\mathbf{u} \equiv (v_1, v_2, u_1, u_2)$

$$s \rightarrow \check{R}(u-v) ; s(v_1, v_2, u_1, u_2) = (u_1, u_2, v_1, v_2)$$

generators of  $\mathfrak{S}_4$  – elementary permutations  $s_1, s_2, s_3$

$$s_1 \mathbf{u} = (v_2, v_1, u_1, u_2) ; s_2 \mathbf{u} = (v_1, u_1, v_2, u_2) ; s_3 \mathbf{u} = (v_1, v_2, u_2, u_1)$$

# Yang-Baxter equation

correspondence  $s_i \rightarrow S_i(\mathbf{u})$  ;  $s_i s_j \rightarrow S_i(s_j \mathbf{u}) S_j(\mathbf{u})$

$$\overbrace{(v_1, v_2, u_1, u_2)}^{S_1} : S_1(\mathbf{u}) L_2(v_1, v_2) = L_2(v_2, v_1) S_1(\mathbf{u})$$

$$\overbrace{(v_1, v_2, u_1, u_2)}^{S_2} : S_2(\mathbf{u}) L_1(u_1, u_2) L_2(v_1, v_2) = L_1(v_2, u_2) L_2(v_1, u_1) S_2(\mathbf{u}) .$$

$$\overbrace{(v_1, v_2, u_1, u_2)}^{S_3} : S_3(\mathbf{u}) L_1(u_1, u_2) = L_1(u_2, u_1) S_3(\mathbf{u})$$

$$S_1(\mathbf{u}) = [i\partial_2]^{v_2-v_1} \quad ; \quad S_2(\mathbf{u}) = [z_{12}]^{u_1-v_2} \quad ; \quad S_3(\mathbf{u}) = [i\partial_1]^{u_2-u_1}$$

defining relations for generators of permutation group  $\mathfrak{S}_4$

quadratic relations are simple

$$s_1 s_1 = \mathbf{1} \rightarrow S_1(s_1 \mathbf{u}) S_1(\mathbf{u}) = [i\partial_2]^{v_2-v_1} \cdot [i\partial_2]^{v_1-v_2} = \mathbf{1}$$

cubic relations reduce to star-triangle relation  $[i\partial]^a \cdot [z]^{a+b} \cdot [i\partial]^b = [z]^b \cdot [i\partial]^{a+b} \cdot [z]^a$  [Isaev]

$$s_1 s_2 s_1 = s_2 s_1 s_2 \rightarrow S_1(s_2 s_1 \mathbf{u}) S_2(s_1 \mathbf{u}) S_1(\mathbf{u}) = S_2(s_1 s_2 \mathbf{u}) S_1(s_2 \mathbf{u}) S_2(\mathbf{u})$$

$$[i\partial_2]^{u_1-v_2} \cdot [z_{12}]^{u_1-v_1} \cdot [i\partial_2]^{v_2-v_1} = [z_{12}]^{v_2-v_1} \cdot [i\partial_2]^{u_1-v_1} \cdot [z_{12}]^{u_1-v_2} .$$



# Yang-Baxter equation

$$s = s_2 s_1 s_3 s_2 \rightarrow \check{R}(u - v) = S_2(s_1 s_3 s_2 \mathbf{u}) S_1(s_3 s_2 \mathbf{u}) S_3(s_2 \mathbf{u}) S_2(\mathbf{u}) =$$

$$= [z_{12}]^{u_2 - v_1} [i\partial_2]^{u_1 - v_1} [i\partial_1]^{u_2 - v_2} [z_{12}]^{u_1 - v_2}$$

useful decomposition on permutations  $u_k \leftrightarrow v_k \quad : \quad \mathfrak{S}_4 \rightarrow \mathfrak{S}_2 \times \mathfrak{S}_2$

$$s = s_2 s_1 s_2 \cdot s_2 s_3 s_2 = r_1 r_2 \rightarrow \check{R}(u - v) = R_1(r_2 \mathbf{u}) R_2(\mathbf{u})$$

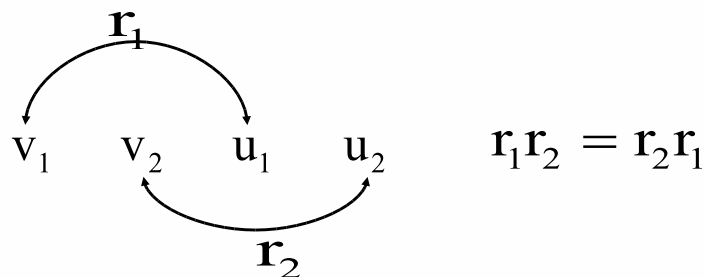
$$r_1 = s_2 s_1 s_2 \rightarrow R_1(\mathbf{u}) = [z_{12}]^{v_2 - v_1} [i\partial_2]^{u_1 - v_1} [z_{12}]^{u_1 - v_2},$$

$$r_2 = s_2 s_3 s_2 \rightarrow R_2(\mathbf{u}) = [z_{12}]^{u_2 - u_1} [i\partial_1]^{u_2 - v_2} [z_{12}]^{u_1 - v_2}.$$

defining equations

$$R_1(\mathbf{u}) L_1(u_1, u_2) L_2(v_1, v_2) = L_1(v_1, u_2) L_2(u_1, v_2) R_1(\mathbf{u})$$

$$R_2(\mathbf{u}) L_1(u_1, u_2) L_2(v_1, v_2) = L_1(u_1, v_2) L_2(v_1, u_2) R_2(\mathbf{u})$$

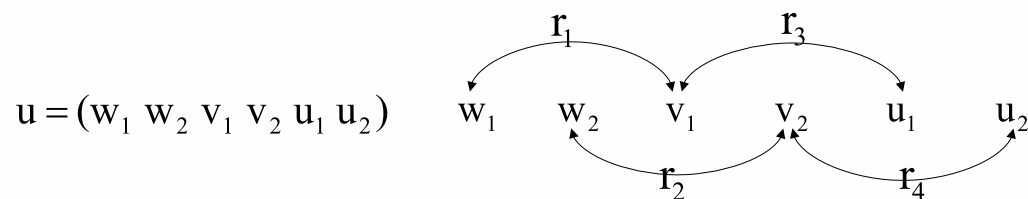
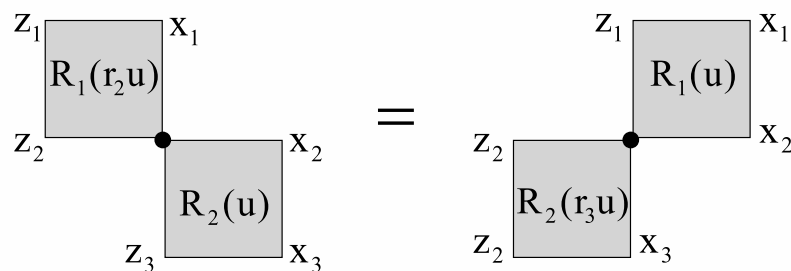
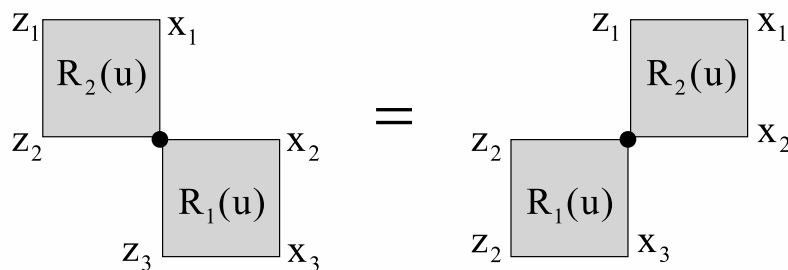


# Yang-Baxter equation

relations in the more complicated case:  $L_1(u_1, u_2) L_2(v_1, v_2) L_3(w_1, w_2)$

$$r_4 r_1 = r_1 r_4 \leftrightarrow R_{12}^2(r_1 \mathbf{u}) R_{23}^1(\mathbf{u}) = R_{23}^1(r_4 \mathbf{u}) R_{12}^2(\mathbf{u})$$

$$r_2 r_3 = r_3 r_2 \leftrightarrow R_{23}^2(r_3 \mathbf{u}) R_{12}^1(\mathbf{u}) = R_{12}^1(r_2 \mathbf{u}) R_{23}^2(\mathbf{u})$$



$$\mathbf{u} = (w_1 \ w_2 \ v_1 \ v_2 \ u_1 \ u_2)$$

# Factorization of general transfer matrix

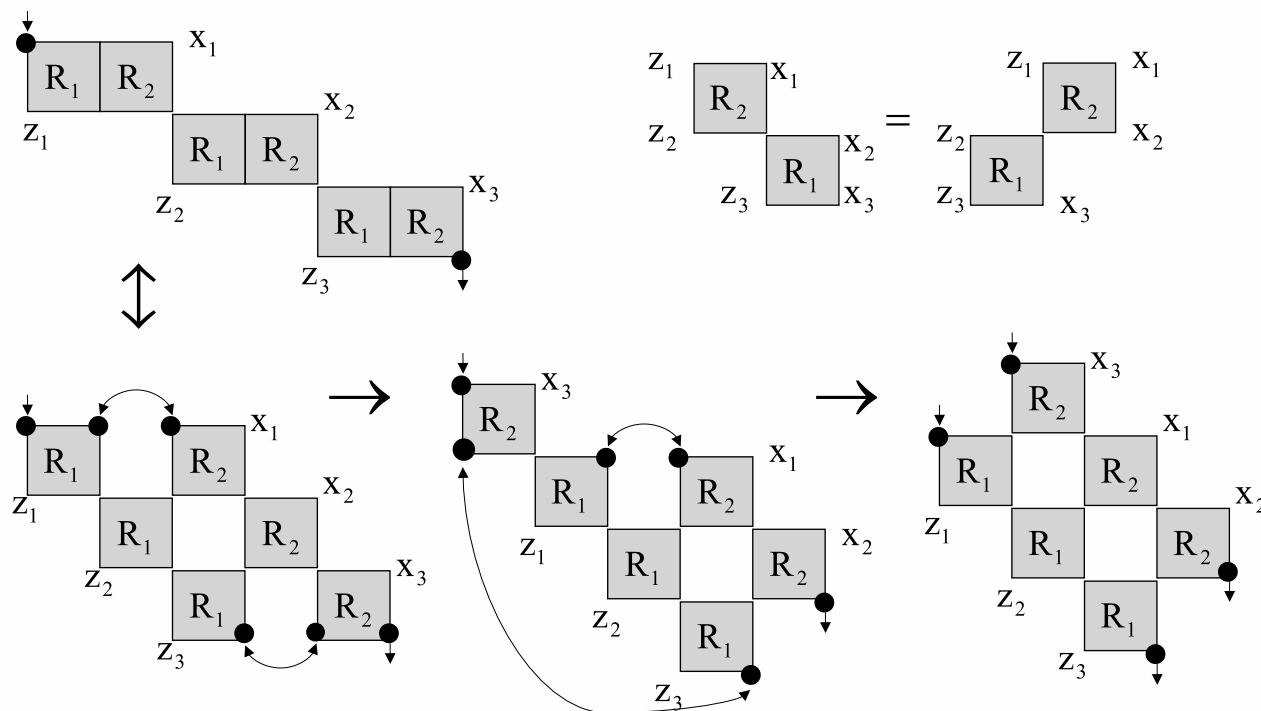
$$\mathbb{T}_{\boldsymbol{\rho}}(u) = \text{tr}_{\mathbb{V}_0} R_{10}(u) R_{20}(u) \dots R_{N0}(u) ; \boldsymbol{\rho} = (\rho_1, \rho_2)$$

commutativity:  $\mathbb{T}_{\boldsymbol{\rho}}(u) \mathbb{T}_{\boldsymbol{\rho}'}(v) = \mathbb{T}_{\boldsymbol{\rho}'}(v) \mathbb{T}_{\boldsymbol{\rho}}(u)$

factorization:  $\mathbb{T}_{\boldsymbol{\rho}}(u) = Q_1(u + \rho_1) \cdot P \cdot Q_2(u + \rho_2) ; P\Phi(z_1, \dots, z_N) = \Phi(z_N, z_1, \dots, z_{N-1})$

$$Q_2(u) = \text{tr}_{\mathbb{V}_0} P_{10} R_{10}^2(u) \cdot P_{20} R_{20}^2(u) \dots P_{N0} R_{N0}^2(u),$$

$$Q_1(u) = \text{tr}_{\mathbb{V}_0} P_{10} R_{10}^1(u) \cdot P_{20} R_{20}^1(u) \dots P_{N0} R_{N0}^1(u)$$



## connection between $Q$ -operators and degeneration points

### connection between $Q$ -operators

$$Q_1(u) = Q_1(u_1, u_2) \quad ; \quad Q_2(v) = Q_2(v_1, v_2)$$

$$Q_2(u_1, u_2) \cdot S = S \cdot Q_1(u_2, u_1) \quad ; \quad S = [i\partial_1]^{u_1-u_2} \dots [i\partial_N]^{u_1-u_2}$$

### degeneration points for $R$ -operators

$$R_1(\mathbf{u})|_{v_1=u_1} = \mathbf{1} \quad ; \quad R_2(\mathbf{u})|_{v_2=u_2} = \mathbf{1}$$

$$R_k(u) = R_k(\mathbf{u})|_{v=0} \rightarrow R_k(u)|_{u=\sigma_k} = \mathbf{1}$$

### degeneration points for $Q$ -operators

$$Q_1(\sigma_1) = P^{-1} \quad ; \quad Q_2(\sigma_2) = P^{-1}$$

### degeneration points for transfer matrix

$$\mathbb{T}_{\boldsymbol{\rho}}(u) = Q_1(u + \rho_1) \cdot P \cdot Q_2(u + \rho_2)$$

$$\boldsymbol{\rho}_2 = (\sigma_1 - u, \sigma_2 + u) \quad ; \quad \boldsymbol{\rho}_1 = (\sigma_1 + u, \sigma_2 - u)$$

$$\mathbb{T}_{\boldsymbol{\rho}_2}(u) = Q_2(2u + \sigma_2) \quad ; \quad \mathbb{T}_{\boldsymbol{\rho}_1}(u) = Q_1(2u + \sigma_1)$$

### commutativity

$$[Q_1(u), Q_1(v)] = [Q_1(u), Q_2(v)] = [Q_2(u), Q_2(v)] = 0$$

## Baxter equation

defining equation for the operator  $R_2(\mathbf{u})$

$$R_2(\mathbf{u}) L_1(u_1, u_2) L_2(v_1, v_2) = L_1(u_1, v_2) L_2(v_1, u_2) R_2(\mathbf{u}) \quad ; \quad [R_2(\mathbf{u}), z_2] = 0$$

$$L_1(u) = Z_1 \cdot \begin{pmatrix} u_1 & -\partial_{z_1} \\ 0 & u_2 \end{pmatrix} \cdot Z_1^{-1} \quad ; \quad L_2(v) = Z_2 \cdot \begin{pmatrix} v_1 & -\partial_{z_2} \\ 0 & v_2 \end{pmatrix} \cdot Z_2^{-1} \quad ; \quad Z_k = \begin{pmatrix} 1 & 0 \\ z_k & 1 \end{pmatrix}$$

$$Z_1^{-1} R_2(\mathbf{u}) L_1(u_1, u_2) Z_2 = \begin{pmatrix} u_1 & -\partial_{z_1} \\ 0 & v_2 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -z_{12} & 1 \end{pmatrix} \begin{pmatrix} v_1 & -\partial_{z_2} \\ 0 & u_2 \end{pmatrix} \cdot R_2(\mathbf{u}) \cdot \begin{pmatrix} v_1 & -\partial_{z_2} \\ 0 & v_2 \end{pmatrix}^{-1}$$

$$Z_1^{-1} R_2(\mathbf{u}) L_1(u_1, u_2) Z_2 = \begin{pmatrix} R_2(\mathbf{u})|_{u_k \rightarrow u_k+1} + v_2 R_2(\mathbf{u}) & -R_2(\mathbf{u}) \partial_{z_1} \\ -v_2 z_{12} R_2(\mathbf{u}) & (u_1 - v_2)(u_2 - v_2) \cdot R_2(\mathbf{u})|_{u_k \rightarrow u_k-1} \end{pmatrix}$$

Baxter equation

$$P_{k0} R_{k0}^2(u) L_k(u_1, u_2) = Z_0 \cdot \begin{pmatrix} P_{k0} R_{k0}^2(u+1) & -P_{k0} R_{k0}^2(u) \partial_k \\ 0 & u_1 u_2 P_{k0} R_{k0}^2(u-1) \end{pmatrix} \cdot Z_0^{-1}$$

$$Q_2(u) = \text{tr}_{V_0} P_{10} R_{10}^2(u) \cdots P_{N0} R_{N0}^2(u) \quad ; \quad t(u) = \text{tr} L_1(u) L_2(u) \cdots L_N(u)$$

$$t(u) Q_2(u) = Q_2(u+1) + (u_1 u_2)^N \cdot Q_2(u-1)$$

## iterative diagonalization of $B(u)$

$$S_- U(\mathbf{z}, \mathbf{x}, p) = p \cdot U(\mathbf{z}, \mathbf{x}, p) ; B(u)U(\mathbf{z}, \mathbf{x}, p) = p \cdot (u - x_1)(u - x_2) \cdots (u - x_{N-1})U(\mathbf{z}, \mathbf{x}, p)$$

defining property of operator  $R_2$

$$\Lambda_N(\mathbf{u}) = R_{12}^2(\mathbf{u})R_{23}^2(\mathbf{u}) \cdots R_{N-1,N}^2(\mathbf{u})$$

$$\Lambda_N(\mathbf{u}) \cdot L_1(u_1, u_2) L_2(u_1, u_2) \cdots L_N(u_1, v_2) = L_1(u_1, v_2) L_2(u_1, u_2) \cdots L_N(u_1, u_2) \cdot \Lambda_N(\mathbf{u})$$

applying to function  $\Psi(z_1 \cdots z_{N-1})$  which does not depend on  $z_N$

$$L_N(u_1, v_2) \rightarrow \begin{pmatrix} u_1 & 0 \\ (u_1 - v_2 + 1) \cdot z_N & v_2 \end{pmatrix} ; L_1(u_1, v_2) = \begin{pmatrix} 1 & 0 \\ z_1 \cdot \frac{u_2 - v_2}{u_2} & \frac{v_2}{u_2} \end{pmatrix} \cdot L_1(u_1, u_2)$$

$$\begin{aligned} \Lambda_N(\mathbf{u}) \cdot \begin{pmatrix} A_{N-1}(u) & B_{N-1}(u) \\ C_{N-1}(u) & D_{N-1}(u) \end{pmatrix} \cdot \Psi(z_1 \cdots z_{N-1}) \cdot \begin{pmatrix} u_1 & 0 \\ (u_1 - v_2 + 1) \cdot z_N & v_2 \end{pmatrix} = \\ = \begin{pmatrix} 1 & 0 \\ z_1 \cdot \frac{u_2 - v_2}{u_2} & \frac{v_2}{u_2} \end{pmatrix} \begin{pmatrix} A_N(u) & B_N(u) \\ C_N(u) & D_N(u) \end{pmatrix} \cdot \Lambda_N(\mathbf{u}) \Psi(z_1 \cdots z_{N-1}) \end{aligned}$$

relation (12)

$$B_N(u) \cdot \Lambda_N(\mathbf{u}) \Psi(z_1 \cdots z_{N-1}) = v_2 \cdot \Lambda_N(\mathbf{u}) \cdot B_{N-1}(u) \Psi(z_1 \cdots z_{N-1})$$

## iterative diagonalization of $B(u)$

usual notations with spectral parameter

$$x = v_2 - u \quad ; \quad \Lambda_N(x) = R_{12}^2(x) R_{23}^2(x) \cdots R_{N-1,N}^2(x)$$

$$B_N(u) \cdot \Lambda_N(x) \Psi(z_1 \cdots z_{N-1}) = (u - x) \cdot \Lambda_N(x) \cdot B_{N-1}(u) \Psi(z_1 \cdots z_{N-1})$$

iteration

$$B_N(u) \cdot \Lambda_N(x_1) \cdots \Lambda_2(x_{N-1}) \Psi(z_1) = (u - x_1) \cdot (u - x_{N-1}) \cdot B_1(u) \cdot \Psi(z_1)$$

$$\Lambda_k(x) = R_{12}^2(x) R_{23}^2(x) \cdots R_{k-1,k}^2(x)$$

one-point operator

$$B_1(u) = -\partial_1 \quad ; \quad B_1(u) e^{-pz_1} = p \cdot e^{-pz_1}$$

answer

$$B(u) U(\mathbf{z}, \mathbf{x}, p) = p \cdot (u - x_1)(u - x_2) \cdots (u - x_{N-1}) U(\mathbf{z}, \mathbf{x}, p)$$

$$U(\mathbf{z}, \mathbf{x}, p) = \Lambda_N(x_1) \cdots \Lambda_2(x_{N-1}) e^{-pz_1}$$

connection with  $Q_2$ -operator

$$Q_2(x) = R_{12}^2(x) R_{23}^2(x) \cdots R_{N-1,N}^2(x) R_{N,0}^2(x) \Big|_{z_0=z_1} \quad ; \quad U(\mathbf{z}, \mathbf{x}, p) = Q_2(x_1) \cdots Q_2(x_{N-1}) e^{-pz_1}$$

## representations of $GL(n, C)$

$$z = \begin{pmatrix} 1 & 0 & 0 & \dots & 0 \\ z_{21} & 1 & 0 & \dots & 0 \\ z_{31} & z_{32} & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ z_{n1} & z_{n2} & z_{n3} & \dots & 1 \end{pmatrix} ; \quad h = \begin{pmatrix} h_{11} & h_{12} & h_{13} & \dots & h_{1n} \\ 0 & h_{22} & h_{23} & \dots & h_{2n} \\ 0 & 0 & h_{33} & \dots & h_{3n} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & h_{n,n} \end{pmatrix}$$

**main construction:** representation with labels  $\sigma = (\sigma_1, \dots, \sigma_n)$  and  $\bar{\sigma} = (\bar{\sigma}_1, \dots, \bar{\sigma}_n)$ ,  $\sigma_k - \bar{\sigma}_k \in \mathbb{Z}$

$$g^{-1} \cdot z = z' \cdot h \quad ; \quad T(g) \Phi(z) = [h_{11}]^{\sigma_1+1} [h_{22}]^{\sigma_2+2} \dots [h_{nn}]^{\sigma_n+n} \cdot \Phi(z')$$

**intertwining operators:** representations  $\sigma$  and  $\sigma_k = s_k \sigma$  are equivalent

$$S_k T^\sigma = T^{\sigma_k} S_k ; \quad k = 1, \dots, n-1$$

$$s_k (\dots \sigma_k, \sigma_{k+1}, \dots) = (\dots \sigma_{k+1}, \sigma_k, \dots)$$

$$S_k = [iD_k]^{\sigma_k, k+1} ; \quad D_k = \frac{\partial}{\partial z_{k+1,k}} + \sum_{m=k+2}^n z_{m,k+1} \frac{\partial}{\partial z_{mk}}$$



## representations of $GL(n, C)$

**L-operator:**  $L(u) = u + \mathbf{e}_{ik} E_{ki}$

explicit formulae for  $GL(3, C)$

$$L(u) = \begin{pmatrix} u + E_{11} & E_{21} & E_{31} \\ E_{12} & u + E_{22} & E_{32} \\ E_{13} & E_{23} & u + E_{33} \end{pmatrix}$$

$$L(u) = \begin{pmatrix} 1 & 0 & 0 \\ z_{21} & 1 & 0 \\ z_{31} & z_{32} & 1 \end{pmatrix} \begin{pmatrix} u - \sigma_1 & -D_1 & -\partial_{z_{31}} \\ 0 & u - \sigma_2 & -D_2 \\ 0 & 0 & u - \sigma_3 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ z_{21} & 1 & 0 \\ z_{31} & z_{32} & 1 \end{pmatrix}^{-1}$$

$$D_1 = \partial_{z_{21}} + z_{32} \partial_{z_{31}} \quad ; \quad D_2 = \partial_{z_{32}}$$

**general structure**

diagonal:  $u - \sigma_1, u - \sigma_2, \dots, u - \sigma_n \leftrightarrow u_1, u_2, \dots, u_n$

next diagonal:  $-D_1, -D_2, \dots, -D_{n-1}$

**intertwining operators  $S_k$  exchange parameters in L-operator**

$$S_k = [iD_k]^{u_{k+1} - u_k} \quad ; \quad S_k L(u_1, \dots, u_k, u_{k+1}, \dots, u_n) = L(u_1, \dots, u_{k+1}, u_k, \dots, u_n) S_k$$

# Yang-Baxter equation

defining relation

$$\check{R}(u-v) L_1(u_1 \cdots u_n) L_2(v_1 \cdots v_n) = L_1(v_1 \cdots v_n) L_2(u_1 \cdots u_n) \check{R}(u-v)$$

permutation group  $\mathfrak{S}_{2n}$  of  $2n$  parameters:  $\mathbf{u} \equiv (v_1, \dots, v_n, u_1, \dots, u_n)$

$$s \leftrightarrow \check{R} \quad ; \quad (v_1, \dots, v_n, u_1, \dots, u_n) \xrightarrow{s} (u_1, \dots, u_n, v_1, \dots, v_n)$$

$$\left( \overbrace{S_1, \dots, S_{n-1}}^{S_1, \dots, S_{n-1}}, \overbrace{S_{n+1}, \dots, S_{2n-1}}^{S_{n+1}, \dots, S_{2n-1}} \right) ; S_k = \begin{cases} \mathbf{1} \otimes S_k, & k = 1, \dots, n-1 \\ S_{k-n} \otimes \mathbf{1}, & k = n+1, \dots, 2n-1 \end{cases}$$

$$S_k L_2(v_1, \dots, v_k, v_{k+1}, \dots, v_n) = L_2(v_1, \dots, v_k, v_{k+1}, \dots, v_n) S_k \quad ; \quad k = 1 \cdots n-1$$

$$S_{n-k} L_1(u_1, \dots, u_k, u_{k+1}, \dots, u_n) = L_1(u_1, \dots, u_k, u_{k+1}, \dots, u_n) S_{n-k} \quad ; \quad k = n+1 \cdots 2n-1$$

$$(v_1, \dots, \overbrace{v_n, u_1, \dots, u_n}^{S_n}).$$

$$S_n L_1(u_1, u_2, \dots, u_n) L_2(v_1, \dots, v_{n-1}, v_n) = L_1(v_n, u_2, \dots, u_n) L_2(v_1, \dots, v_{n-1}, u_1) S_n.$$

$$S_n = [z_2^{-1} z_1]_{n1}^{u_1 - v_n} \quad ; \quad z_2^{-1} z_1 = \begin{pmatrix} 1 & 0 \\ z_2 & 1 \end{pmatrix}^{-1} \begin{pmatrix} 1 & 0 \\ z_1 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ z_1 - z_2 & 1 \end{pmatrix}$$

# Yang-Baxter equation

defining relations for generators

$$\mathbb{S}_k(s_k \mathbf{u}) \mathbb{S}_k(\mathbf{u}) = \mathbf{1}$$

$$\mathbb{S}_i(s_k \mathbf{u}) \mathbb{S}_k(\mathbf{u}) = \mathbb{S}_k(s_i \mathbf{u}) \mathbb{S}_i(\mathbf{u}) ; |i - k| > 1$$

$$\mathbb{S}_k(s_{k+1} s_k \mathbf{u}) \mathbb{S}_{k+1}(s_k \mathbf{u}) \mathbb{S}_k(\mathbf{u}) = \mathbb{S}_{k+1}(s_k s_{k+1} \mathbf{u}) \mathbb{S}_k(s_{k+1} \mathbf{u}) \mathbb{S}_{k+1}(\mathbf{u})$$

useful decomposition on permutations  $u_k \leftrightarrow v_k$ :  $\mathfrak{S}_{2n} \rightarrow \mathfrak{S}_2 \times \cdots \times \mathfrak{S}_2$

$$(u_1, \dots, u_k, \dots, v_k, \dots, v_n) \xrightarrow{r_k} (u_1, \dots, v_k, \dots, u_k, \dots, v_n)$$

$$R_k L_1(u_1 \dots u_k \dots u_n) L_2(v_1 \dots v_k \dots v_n) = L_1(u_1 \dots v_k \dots u_n) L_2(v_1 \dots u_k \dots v_n) R_k$$

$$R_k = (\mathbb{S}_{n+k-1} \dots \mathbb{S}_{n+1}) (\mathbb{S}_k \dots \mathbb{S}_{n-1}) \mathbb{S}_n (\mathbb{S}_{n-1} \dots \mathbb{S}_k) (\mathbb{S}_{n+1} \dots \mathbb{S}_{n+k-1})$$

factorization of R-operator and relations needed for factorization of transfer-matrix

$$\check{R}(\mathbf{u}) = R_1(r_2 \cdots r_n \mathbf{u}) \dots R_{n-1}(r_n \mathbf{u}) R_n(\mathbf{u})$$

$$R_{12}^k(\mathbf{u}) R_{23}^i(\mathbf{u}) = R_{23}^i(\mathbf{u}) R_{12}^k(\mathbf{u}) \quad ; \quad k > i$$

# Factorization of transfer matrix and degeneration points

## factorization of transfer-matrix

$$\mathbb{T}_{\boldsymbol{\rho}}(u) = \text{tr}_{\mathbb{V}_0} R_{10}(u) R_{20}(u) \dots R_{N0}(u) ; \boldsymbol{\rho} = (\rho_1, \dots, \rho_n)$$

$$\mathbb{T}_{\boldsymbol{\rho}}(u) = Q_1(u + \rho_1) \cdot P \cdot Q_2(u + \rho_2) \dots P \cdot Q_n(u + \rho_n)$$

$$Q_k(u) = \text{tr}_{\mathbb{V}_0} P_{10} R_{10}^k(u) \cdot P_{20} R_{20}^k(u) \dots P_{N0} R_{N0}^k(u)$$

## degeneration points for R-operators

$$R_k(\mathbf{u})|_{v_k=u_k} = \mathbf{1}$$

$$R_k(u) = R_k(\mathbf{u})|_{v=0} \rightarrow R_k(u)|_{u=\sigma_k} = \mathbf{1}$$

## degeneration points for Q-operators

$$Q_k(\sigma_k) = P^{-1}$$

## degeneration points for transfer matrix

$$\boldsymbol{\rho}_k = (\sigma_1 - u, \sigma_2 - u, \dots, \sigma_k + u \cdot (n - 1), \dots, \sigma_n - u)$$

$$\mathbb{T}_{\boldsymbol{\rho}_k}(u) = Q_k(nu + \sigma_k)$$

## commutativity

$$\mathbb{T}_{\boldsymbol{\rho}}(u) \mathbb{T}_{\boldsymbol{\rho}'}(v) = \mathbb{T}_{\boldsymbol{\rho}'}(v) \mathbb{T}_{\boldsymbol{\rho}}(u) \rightarrow [Q_i(u), Q_k(v)] = 0$$

# Baxter equation

defining relation for  $R_n$

$$R_{k0}^n(\mathbf{u}) L_k(u_1, \dots, u_n) L_0(v_1, \dots, v_n) = L_k(u_1, \dots, v_n) L_0(v_1, \dots, u_n) R_{k0}^n(\mathbf{u})$$

at the point  $v_n = 0$  can be rewritten in the form

$$P_{k0} R_{k0}^n(u) L_k(u_1, \dots, u_n) = Z_0 \cdot \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1n} \\ a_{21} & a_{22} & \dots & a_{2n} \\ \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & \dots & u_1 \cdots u_n \cdot P_{k0} R_{k0}^n(u-1) \end{pmatrix} \cdot Z_0^{-1}$$

transfer matrices

$$\mathbb{T}(u) = L_1(u) \cdots L_N(u)$$

$$t_k(u) = \text{tr}_k \mathbb{A}_k \mathbb{T}_1(u) \mathbb{T}_2(u+1) \cdots \mathbb{T}_k(u+k-1)$$

$\mathbb{A}_k$  - anti-symmetrization of  $k$  indices and trace in  $\otimes^k \mathbb{C}^n$

explicit formulae

$$(\mathbb{A}_k)_{a_1 a_2 \dots a_k}^{b_1 b_2 \dots b_k} = \frac{1}{k!} \sum_p (-1)^{\text{sign}(p)} \delta_{p(a_1)}^{b_1} \delta_{p(a_2)}^{b_2} \cdots \delta_{p(a_k)}^{b_k}$$

$$t_k(u) = \sum_{a_1 < a_2 < \dots < a_k} \sum_p (-1)^{\text{sign}(p)} \mathbb{T}_{p(a_1) a_1}(u) \mathbb{T}_{p(a_2) a_2}(u+1) \cdots \mathbb{T}_{p(a_k) a_k}(u+k-1)$$

# Baxter equation

## examples of Baxter equation

$$A(u) = (u_1 \cdots u_n)^N ; \quad t_n(u) = A(u + n - 1) ; \quad Q(u) = Q_n(u)$$

✓  $n = 2$

$$A(u + 1) \cdot Q(u) t_1(u) - Q(u + 1) t_2(u) = A(u) A(u + 1) \cdot Q(u - 1)$$

$$Q(u) t_1(u) - Q(u + 1) = u_1^N u_2^N \cdot Q(u - 1)$$

✓  $n = 3$

$$A(u+1)A(u+2) \cdot Q(u)t_1(u) - A(u+2) \cdot Q(u+1)t_2(u) + Q(u+2)t_3(u) = A(u)A(u+1)A(u+2) \cdot Q(u-1)$$

$$A(u + 1) \cdot Q(u) t_1(u) - Q(u + 1) t_2(u) + Q(u + 2) = A(u) A(u + 1) \cdot Q(u - 1)$$

in the general case all can be packed in one formula

$$\text{tr}_n \mathbb{A}_n \left[ e^{\partial_u} \mathbb{T}_1(u) - A(u + n) \right] \left[ e^{\partial_u} \mathbb{T}_2(u) - A(u + n - 1) \right] \left[ e^{\partial_u} \mathbb{T}_n(u) - A(u + 1) \right] \cdot Q_n(u) = 0$$

$\mathbb{A}_n$  - anti-symmetrization of  $n$  indices and trace in  $\otimes^n \mathbb{C}^n$

[conjectured by Chervov-Talalaev]

## fusion relations

- ✓ For integer points:  $\rho_k - \rho_{k+1} \in \mathbb{N}$  the representation with label  $\rho$  is not irreducible: in infinite-dimensional space  $\mathbb{V}_\rho$  there exists finite-dimensional invariant subspace.
- ✓ Transfer-matrix with this finite-dimensional auxiliary space  $t_\rho(u)$  can be expressed in terms of  $\mathbb{T}_\rho(u)$

$$t_\rho(u) = \sum_{\mathbf{p}} (-)^{\text{sign}(\mathbf{p})} \cdot \mathbb{T}_{\mathbf{p}\rho}(u) \quad (1)$$

where  $\mathbf{p}\rho = (\rho_{k_1}, \rho_{k_2} \cdots \rho_{k_n})$  is permutation of  $(\rho_1, \rho_2 \cdots \rho_n)$ . The sum is over all permutations.

- ✓ The factorization of  $\mathbb{T}_\rho(u)$  into product of Q-operators leads to the following determinant representation for  $t_\rho(u)$

$$t_\rho(u) = P^{n-1} \cdot \begin{vmatrix} Q_1(u + \rho_1) & Q_2(u + \rho_1) & \cdots & Q_n(u + \rho_1) \\ Q_1(u + \rho_2) & Q_2(u + \rho_2) & \cdots & Q_n(u + \rho_2) \\ \cdots & \cdots & \cdots & \cdots \\ Q_1(u + \rho_n) & Q_2(u + \rho_n) & \cdots & Q_n(u + \rho_n) \end{vmatrix}$$

## conclusions

- ✓ simplest building block – operator  $R_n(\mathbf{u})$  which interchanges  $u_n \leftrightarrow v_n$

$$R_n(\mathbf{u}) L_1(u_1, \dots, u_n) L_2(v_1, \dots, v_n) = L_1(u_1, \dots, v_n) L_2(v_1, \dots, u_n) R_n(\mathbf{u})$$

- ✓  $R_n(\mathbf{u}) \leftrightarrow Q\text{-operator} \leftrightarrow U(\mathbf{z}, \mathbf{x})$  - transition operator to the SOV-representation

$$\Psi(\mathbf{z}) = \int U(\mathbf{z}, \mathbf{x}) Q(\mathbf{x}) \mu(\mathbf{x}) d^N \mathbf{x}$$

- ✓ The last step is still missing for  $SL(n, C)$ .